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# $\mathbf{E}_{8}$ lattice with icosians and $\mathbf{Z}_{5}$ symmetry 

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#### Abstract

A simple method for the construction of the $\mathrm{E}_{8}$ root system with icosians is suggested. It is confronted with the root system of $E_{8}$ obtained by integral octonions. Embeddings of the maximal subgroups $\mathrm{SU}(5) \times \mathrm{SU}(5)$ and $\mathrm{SU}(3) \times \mathrm{E}_{6}$ in $\mathrm{E}_{8}$ with respective five-fold and three-fold symmetries are discussed.


## 1. Introduction

Most general string theories can be obtained using self-dual lattices where the compactified degrees of freedom are described by simple conformal field theories [1]. Among these theories the heterotic string [2] constructed in ten-dimensional spacetime displays most attractive features with a possibility of obtaining the standard model with orbifold compactification [3]. With regard to these facts the root system of $E_{8}$ plays an essential role. In a recent paper [4] we have constructed the root system of $\mathrm{E}_{8}$ with integral octonions [5] and investigated its algebraic structure. The method which we have employed is based on the idea that the octonions are obtained by pairing two quaternion $\bar{s}$. For this purpose we first obtained the root system of $\mathrm{F}_{4}$ and then combined two such systems to construct the octonionic description of the $E_{8}$ lattice, where one set of $\mathrm{F}_{4}$ roots is multiplied by an imaginary unit $e_{7}$ and added to the other.

In this paper we construct the $\mathrm{E}_{8}$ lattice with icosians. Icosian is a generic name for the 120 quaternionic elements $(q)$ of the binary icosahedral group [6] which we will discuss in §2. We follow the same method as [4], i.e. we combine two sets of quaternionic roots of $\mathrm{F}_{4}$, multiplying one set by $\sigma=\frac{1}{2}(1-\sqrt{5})$ and add it to the other set, which leads to 240 non-zero roots $q, \sigma q$ of $\mathrm{E}_{8}$. We compare the $\mathrm{E}_{8}$ roots of icosians with those of octonions and find the relations between them. In § 3 we decompose the roots of $\mathrm{E}_{8}$ with respect to one of its maximal subgroups $\mathrm{SU}(5) \times \mathrm{SU}(5)^{\prime}$ where a five-fold symmetry of icosians plays a dominant role. In § 4 we concentrate on the three-fold symmetry of icosians by branching $\mathrm{E}_{8}$ with respect to its maximal subgroup $\operatorname{SU}(3) \times E_{6}$. Section 5 consists of the discussions and suggestions as to how this method can be generalised for the construction of the $E_{8} \times E_{8}^{\prime}$ lattice with the inclusion of octonions and the Leech lattice [7]. A preliminary version of this work has been published [8].

[^0]
## 2. Icosians as roots of $\mathbf{E}_{8}$

A discrete subgroup of $\mathrm{SO}(3)$ of order 60 is called the icosahedral group, which is the symmetry group of the icosahedron. It is isomorphic to the group of even permutations $\mathrm{A}_{5}$ of five letters. Its double cover $2 \mathrm{~A}_{5}$ of 120 elements, called the binary icosahedral group, can be represented by quaternions or equivalently $2 \times 2$ unitary matrices of determinant one. The 120 elements of the binary icosahedral group can be generated by the elements [6]

$$
\begin{equation*}
A=\frac{1}{2}\left(\tau-\sigma e_{1}+e_{3}\right) \quad B=\frac{1}{2}\left(1-\sigma e_{2}+\tau e_{3}\right) \tag{1}
\end{equation*}
$$

satisfying the relations

$$
\begin{equation*}
A^{5}=B^{3}=C^{2}=A B C=-1 \tag{2}
\end{equation*}
$$

with $C=e_{3}$. Here $\tau$ and $\sigma$ are defined by

$$
\begin{array}{lll}
\tau=\frac{1}{2}(1+\sqrt{5}) & \sigma=\frac{1}{2}(1-\sqrt{5})  \tag{3}\\
\tau+\sigma=1 & \tau^{2}=\tau+1 & \sigma^{2}=\sigma+1 . \\
\tau \sigma=-1
\end{array}
$$

and $e_{1}, e_{2}$ and $e_{3}$ are the quaternionic imaginary units satisfying

$$
\begin{equation*}
e_{i} e_{j}=-\delta_{i j}+\varepsilon_{i j k} e_{k} \quad \bar{e}_{i}=-e_{i} \quad i, j, k=1,2,3 \tag{4}
\end{equation*}
$$

where $\delta_{i j}$ and $\varepsilon_{i j k}$ are the usual Kronecker and Levi-Civita symbols respectively. In mathematical literature a short notation $\langle p, q, r\rangle$ is used to denote the groups generated by

$$
\begin{equation*}
A^{p}=B^{q}=C^{r}=A B C=-1 \tag{5}
\end{equation*}
$$

Groups generated by quaternions fall into four classes called the quaternion group $\langle 2,2,2\rangle$ of order 8 , the binary tetrahedral group $\langle 3,3,2\rangle$ of order 24 , the binary octahedral group $\langle 4,3,2\rangle$ of order 48 and finally the binary icosahedral group $\langle 5,3,2\rangle$ of order 120. In [4] we have shown the relations of $\langle 3,3,2\rangle$ and $\langle 4,3,2\rangle$ with the quaternionic root systems of $\mathrm{SO}(8)$ and $\mathrm{F}_{4}$ respectively. An explicit form for the elements of $\langle 5,3,2\rangle$ can be calculated using (1) and is written as follows:
$\pm 1 \quad \pm e_{1} \quad \pm e_{2} \quad \pm e_{3} \quad \frac{1}{2}\left( \pm 1 \pm e_{1} \pm e_{2} \pm e_{3}\right)$
$\frac{1}{2}\left( \pm \tau \pm e_{1} \pm \sigma e_{2}\right) \quad \frac{1}{2}\left( \pm 1 \pm \tau e_{1} \pm \sigma e_{3}\right) \quad \frac{1}{2}\left( \pm \sigma \pm \tau e_{2} \pm e_{3}\right) \quad \frac{1}{2}\left( \pm \sigma e_{1} \pm e_{2} \pm \tau e_{3}\right)$
$\frac{1}{2}\left( \pm \tau \pm e_{2} \pm \sigma e_{3}\right) \quad \frac{1}{2}\left( \pm 1 \pm \sigma e_{1} \pm \tau e_{2}\right) \quad \frac{1}{2}\left( \pm \sigma \pm e_{1} \pm \tau e_{3}\right) \quad \frac{1}{2}\left( \pm \tau e_{1} \pm \sigma e_{2} \pm e_{3}\right)$
$\frac{1}{2}\left( \pm \tau \pm \sigma e_{1} \pm e_{3}\right) \quad \frac{1}{2}\left( \pm 1 \pm \sigma e_{2} \pm \tau e_{3}\right) \quad \frac{1}{2}\left( \pm \sigma \pm \tau e_{1} \pm e_{2}\right) \quad \frac{1}{2}\left( \pm e_{1} \pm \tau e_{2} \pm \sigma e_{3}\right)$.
The 24 integral quaternions in ( $6 a$ ) (Hurwitz integers) [9] are the elements of $\langle 3,3,2\rangle$, a subgroup of $\langle 5,3,2\rangle .\langle 4,3,2\rangle$ is not a subgroup of $\langle 5,3,2\rangle$. Notice that ( $6 c$ ) and ( $6 d$ ) follow from ( $6 b$ ) by a cyclic permutation of $e_{1}, e_{2}$ and $e_{3}$. If we denote by $q$ any element of $\langle 5,3,2\rangle$ then there are 30 elements satisfying $q^{2}=-1,40$ elements with $q^{3}= \pm 1$ (half with $q^{3}=+1$ ) and 48 elements with $q^{5}= \pm 1$ (half with $q^{5}=+1$ ). As we shall discuss in $\S \S 3$ and 4 , these features of icosians are appropriate for the decompositions of $E_{8}$ with respect to its maximal subgroups $S U(3) \times E_{6}$ and $S U(5) \times S U(5)^{\prime}$.

Wilson [10], as well as Conway and Sloane [11] have proposed that the $\mathrm{E}_{8}$ lattice can be described by 120 icosians $q$ and their multiples with $\sigma, \sigma q$. To ensure the 'correct angles' between the $\mathrm{E}_{8}$ roots, they suggested a 'reduced' scalar product. Let $p$ and $q$ be two quaternions. The usual scalar product is defined by

$$
\begin{equation*}
(p, q)=\frac{1}{2}(\tilde{p} q+\bar{q} p) \tag{7}
\end{equation*}
$$

With this definition, the scalar products of icosians (6a)-(6d) will take the values $a+b \sigma$, where $a$ and $b$ are $0, \pm \frac{1}{2}, \pm 1$. The 'reduced' scalar product is defined by the mapping $a+b \sigma \rightarrow a$. This new definition of the scalar product also leads to a construction of the Leech lattice in terms of icosians [12]. Therefore, multiplying the elements in ( $6 a$ ) - $6 d$ ) by $\sigma$ will help us to write the complete root system of $\mathrm{E}_{8}$ explicitly. However, in this section, without referring to the defining relation of (2) we give an alternative construction of the roots of $E_{8}$ with icosians. We follow a method similar to that suggested in [4]. We notice that the 'reduced' scalar product allows us to treat $\sigma, \sigma e_{1}, \sigma e_{2}, \sigma e_{3}$ as new orthogonal units, independent of the quaternionic units $1, e_{1}$, $e_{2}$ and $e_{3}$. Thus, by this trick, we enlarge the four-dimensional Euclidean space of quaternions to eight-dimensional Euclidean space, a necessary step towards the construction of the $\mathrm{E}_{8}$ lattice. This procedure immediately suggests that there must exist a natural correspondence between $\sigma, \sigma e_{1}, \sigma e_{2}, \sigma e_{3}$ and the octonionic units $e_{4}, e_{5}, e_{6}$ and $e_{7}$. This point will be clarified in what follows.

We briefly recall the octonionic construction of $E_{8}$ in [4]. The root system of the exceptional group $F_{4}$ can be written in terms of quaternions as follows

| $A_{0}$ | $A_{1}$ | $A_{2}$ |  |
| :--- | :--- | :--- | :--- |
| $24: \pm 1, \pm e_{1}, \pm e_{2}, \pm e_{3}$ <br> $\frac{1}{2}\left( \pm 1 \pm e_{1} \pm e_{2} \pm e_{3}\right)$ |  | $8_{\mathrm{v}}: \frac{1}{2}\left( \pm 1 \pm e_{1}\right)$ <br> $\frac{1}{2}\left( \pm e_{2} \pm e_{3}\right)$ | $8_{\mathrm{c}}: \frac{1}{2}\left( \pm 1 \pm e_{2}\right)$ <br> $\frac{1}{2}\left( \pm e_{3} \pm e_{1}\right)$ |

Here, provided they are multiplied by $\sqrt{2}, A_{0}$ represents the long roots of $F_{4}$ (the root system of $\mathrm{SO}(8)$ ) and $A_{1}, A_{2}$ and $A_{3}$ constitute the short roots of $F_{4}$. By pairing two sets of (8) in a delicate manner we obtain the octonionic representation of $\mathrm{E}_{8}$ roots:

$$
\begin{align*}
& {\left[A_{0}, 0\right]=A_{0} \quad\left[0, A_{0}\right]=e_{7} A_{0} \quad\left[A_{1}, A_{1}\right]=A_{1}+e_{7} A_{1}}  \tag{9}\\
& {\left[A_{2}, A_{3}\right]=A_{2}+e_{7} A_{3} \quad\left[A_{3}, A_{2}\right]=A_{3}+e_{7} A_{2}}
\end{align*}
$$

where we introduce the octonionic units $e_{7}^{2}=-1, \bar{e}_{7}=-e_{7}, e_{4}=e_{7} e_{1}, e_{5}=e_{7} e_{2}$ and $e_{6}=e_{7} e_{3}$. In fact, this is an explicit realisation of the $\mathrm{E}_{8}$ lattice with octonions by pairing two $\mathrm{F}_{4}$ lattices suggested by Goddard et al [13] where $\mathrm{E}_{8}$ sits at one corner of a magic square. A little effort shows that $E_{8}$ can also be constructed by pairing two $\mathrm{F}_{4}$ roots similarly to (9) but with slightly different partners and replacing $e_{7}$ by $\sigma$. Indeed, one can easily check that the following pairings:

$$
\begin{align*}
& \left(A_{0}, 0\right)=A_{0} \quad\left(0, A_{0}\right)=\sigma A_{0} \quad\left(A_{1}, A_{2}\right)=A_{1}+\sigma A_{2}  \tag{10}\\
& \left(A_{2}, A_{3}\right)=A_{2}+\sigma A_{3} \quad\left(A_{3}, A_{1}\right)=A_{3}+\sigma A_{1}
\end{align*}
$$

not only reproduce the 120 elements in $(6 a)-(6 d)$, denoted by $q$, but also yield the additional roots $\sigma q$ of $\mathrm{E}_{8}$. The differences between two pairings can be contrasted by comparing (9) and (10). This comparison suggests that a correspondence between the octonionic roots and the icosians can be obtained in the following form

$$
\begin{array}{rlrl}
e_{7} A_{0} \leftrightarrow \sigma A_{0} & & e_{7} \leftrightarrow \sigma \\
e_{7} A_{1} \leftrightarrow \sigma A_{2} & &  \tag{11}\\
e_{7} A_{2} \leftrightarrow \sigma A_{1} \\
e_{7} A_{3} \leftrightarrow \sigma A_{3} & \rightarrow & e_{7} e_{1}=e_{4} \leftrightarrow \sigma e_{2} \\
e_{7} e_{2}=e_{5} \leftrightarrow \sigma e_{1} \\
e_{7} e_{3}=e_{6} \leftrightarrow \sigma e_{3} .
\end{array}
$$

With the obvious mapping $1 \leftrightarrow 1, e_{1} \leftrightarrow e_{1}, e_{2} \leftrightarrow e_{2}, e_{3} \leftrightarrow e_{3}$, one can easily transform one system of roots of $\mathrm{E}_{8}$ into another. This transformation can also be used for the
octonionic construction of the Leech lattice, which has been already described by icosians [10-12].

Before we end this section, let us remark on the following facts. There exists an alternative representation of the $\mathrm{E}_{8}$ lattice with icosians. Instead of starting with the pair $\left(A_{1}, A_{2}\right)$ in (10), had we started with $\left(A_{1}, A_{3}\right)$ we would have obtained the following set of icosians:

$$
\begin{align*}
& \left(A_{0}, 0\right)=A_{0} \quad\left(0, A_{0}\right)=\sigma A_{0} \quad\left(A_{1}, A_{3}\right)=A_{1}+\sigma A_{3} \\
& \left(A_{3}, A_{2}\right)=A_{3}+\sigma A_{2} \quad\left(A_{2}, A_{1}\right)=A_{2}+\sigma A_{1} . \tag{12}
\end{align*}
$$

The 120 -element subset in (12), which constitutes the binary icosahedral group, is completely independent of ( $6 a$ )-( $6 d$ ) and can be generated by $A=\frac{1}{2}\left(\tau+\sigma e_{2}+e_{3}\right)$ and $B=\frac{1}{2}\left(1-\sigma e_{1}+\tau e_{3}\right)$. Equation (12) is obtained from (10) by a redefinition of the quaternionic units $e_{1} \rightarrow-e_{2}, e_{2} \rightarrow e_{1}, e_{3} \rightarrow e_{3}$ corresponding to a rotation of $\pi / 2$ around the $e_{3}$ axis in the clockwise direction, which can be obtained by the action of an element of the octahedral group. Since the octohedral or binary octohedral group is not a subgroup of $\langle 5,3,2\rangle$ the new set of icosians are expected to be different from the former. The elements of the binary icosahedral group used in most of the mathematical literature are those which can be obtained from (12). If one compares (12) and (9), the correspondence in this case between octonions and icosians can be obtained from the mapping
$\sigma \leftrightarrow e_{7} \quad \sigma e_{1} \leftrightarrow e_{6}=e_{7} e_{3} \quad \sigma e_{2} \leftrightarrow e_{5}=e_{7} e_{2} \quad \sigma e_{3} \leftrightarrow e_{4}=e_{7} e_{1}$.
In the appendix to [4] we have given seven different constructions of the $\mathrm{E}_{8}$ lattice with octonions similar to (9). Indeed, with the quaternionic units $e_{1}, e_{2}$ and $e_{3}$ one can also construct the following two independent octonionic root systems of $\mathrm{E}_{8}$ :

$$
\begin{align*}
& {\left[A_{0}, 0\right]=A_{0} \quad\left[0, A_{0}\right]=e_{7} A_{0} \quad\left[A_{2}, A_{2}\right]=A_{2}+e_{7} A_{2}}  \tag{14}\\
& {\left[A_{3}, A_{1}\right]=A_{3}+e_{7} A_{1} \quad\left[A_{1}, A_{3}\right]=A_{1}+e_{7} A_{3}}
\end{align*}
$$

and

$$
\begin{align*}
& {\left[A_{0}, 0\right]=A_{0} \quad\left[0, A_{0}\right]=e_{7} A_{0} \quad\left[A_{3}, A_{3}\right]=A_{3}+e_{7} A_{3}} \\
& {\left[A_{1}, A_{2}\right]=A_{1}+e_{7} A_{2} \quad\left[A_{2}, A_{1}\right]=A_{2}+e_{7} A_{1} .} \tag{15}
\end{align*}
$$

Relations among these octonionic constructions and those in (10) and (12) can be found in a similar manner which we have illustrated. Although we shall say more in concluding remarks as regards the algebraic aspects of the $\mathrm{E}_{8}$ roots with octonions contrasted with icosians, some of their properties should be mentioned here since we will use them in later sections. Let $P$ represent an octonionic root of $\mathrm{E}_{8}$. They satisfy either conditions $q^{3}= \pm 1, q^{2}= \pm 1$. This feature of the octonionic roots can be used for the constructions of orbifolds with $\mathrm{Z}_{6}, \mathrm{Z}_{4}, \mathrm{Z}_{3}$ and $\mathrm{Z}_{2}$ symmetries embeddable in $\mathrm{E}_{8}$ [14]. In the case of icosianic roots we also have the roots satisfying $q^{5}= \pm 1$ in addition to the former ones, so that it allows constructions of orbifolds with $\mathrm{Z}_{10}$ and $Z_{5}$ symmetries.

## 3. Embedding of $\operatorname{SU}(5) \times \operatorname{SU}(5)^{\prime}$ in $\mathbf{E}_{8}$ with $\mathbf{Z}_{5}$ symmetry

Hereafter we will deal with the icosian description of the $\mathrm{E}_{8}$ lattice given by (10). Any element $q$ in $(6 b)-(6 d)$ which has the scalar part $-\frac{1}{2} \tau$ or $-\frac{1}{2} \sigma$ has the property $q^{5}=1$.

Therefore 24 elements with this property can be classified into six groups. This classification can be done at will depending on the choices of six elements; to begin with let us choose a root $R_{1}=\frac{1}{2}\left(-\tau+e_{1}+\sigma e_{2}\right)$. A simple calculation shows that its powers are given by

$$
\begin{array}{ll}
R_{1}=\frac{1}{2}\left(-\tau+e_{1}+\sigma e_{2}\right) \quad & R_{1}^{2}=\frac{1}{2}\left(-\sigma-\tau e_{1}+e_{2}\right) \\
R_{1}^{3}=\bar{R}_{1}^{2}=\frac{1}{2}\left(-\sigma+\tau e_{1}-e_{2}\right) \quad & R_{1}^{4}=\bar{R}_{1}=\frac{1}{2}\left(-\tau-e_{1}-\sigma e_{2}\right)  \tag{16}\\
R_{1}^{5}=1 & R_{1}+R_{1}^{2}+R_{1}^{3}+R_{1}^{4}+R_{1}^{5}=0 .
\end{array}
$$

One can immediately check that with the use of the 'reduced' scalar product these roots can be used for the description of an extended Coxeter-Dynkin diagram of $\operatorname{SU}(5)$ (figure 1). The 20 non-zero roots of $\mathrm{SU}(5)$ are then given by

$$
\begin{align*}
& \pm R_{1}, \quad \pm R_{1}^{2}, \quad \pm R_{1}^{3}, \quad \pm R_{1}^{4}, \quad \pm R_{1}^{5}= \pm 1  \tag{17}\\
& \pm \sigma R_{1}, \quad \pm \sigma R_{1}^{2}, \quad \pm \sigma R_{1}^{3}, \quad \pm \sigma R_{1}^{4}, \quad \pm \sigma R_{1}^{5}= \pm \sigma .
\end{align*}
$$

It is clear that this set of $\mathrm{SU}(5)$ roots is left invariant under a repeated left or right multiplication of $R_{1}$, leading to a $Z_{5}$ symmetry. An orthogonal set of roots to the roots in (17) can be generated by repeated application of $R_{1}$ on the roots $\pm e_{3}$ and $\pm \sigma e_{3}$. Thus we obtain the roots of another $\mathrm{SU}(5)^{\prime}$, orthogonal to $\mathrm{SU}(5)$, given by

$$
\begin{align*}
& \pm R_{1} e_{3}, \quad \pm R_{1}^{2} e_{3}, \quad \pm R_{1}^{3} e_{3}, \quad \pm R_{1}^{4} e_{3}, \quad \pm e_{3}, \\
& \pm \sigma R_{1} e_{3}, \quad \pm \sigma R_{1}^{2} e_{3}, \quad \pm \sigma R_{1}^{3} e_{3}, \quad \pm \sigma R_{1}^{4} e_{3}, \quad \pm \sigma e_{3} . \tag{18}
\end{align*}
$$

Hence (17) and (18) display the roots of the maximal subgroup $\mathrm{SU}(5) \times \mathrm{SU}(5)^{\prime}$ of $\mathrm{E}_{8}$ with an obvious $Z_{5}$ symmetry. The remaining roots belong to the coset space $\mathrm{E}_{8} / \mathrm{SU}(5) \times$ $\operatorname{SU}(5)^{\prime}$, which transform as $\left(\underline{5}, \underline{10^{*}}\right)+\left(\underline{5}^{*}, \underline{10}\right)$ and $(\underline{10}, \underline{5})+\left(\underline{10}, \underline{5}^{*}\right)$. The 200 roots of the coset space can be written in a compact form by defining
$A_{1}=\frac{1}{2}\left(-1-e_{1}-e_{2}+e_{3}\right) \quad A_{2}=\frac{1}{2}\left(-\sigma-\tau e_{2}+e_{3}\right) \quad A_{3}=\frac{1}{2}\left(-\tau e_{1}+\sigma e_{2}+e_{3}\right)$
$A_{4}=\frac{1}{2}\left(\sigma-\tau e_{2}+e_{3}\right)=-\bar{A}_{2} \quad A_{5}=\frac{1}{2}\left(1-e_{1}-e_{2}+e_{3}\right)=-\bar{A}_{1}$
$B_{1}=\frac{1}{2}\left(\tau-e_{2}+\sigma e_{3}\right) \quad B_{2}=\frac{1}{2}\left(-1-\tau e_{1}+\sigma e_{3}\right) \quad B_{3}=\frac{1}{2}\left(e_{1}-\tau e_{2}+\sigma e_{3}\right)$
$B_{4}=\frac{1}{2}\left(1-\tau e_{1}+\sigma e_{3}\right)=-\bar{B}_{2} \quad B_{5}=\frac{1}{2}\left(-\tau-e_{2}+\sigma e_{3}\right)=-\bar{B}_{1}$.
It can be shown that the roots of the representation ( $5,10^{*}$ ) are given by the elements

$$
\begin{equation*}
R_{1}^{n} A_{m}, \quad \sigma R_{1}^{n} B_{m} \quad n, m=1,2,3,4,5 . \tag{20}
\end{equation*}
$$



Figure 1. Extended Coxeter-Dynkin diagram of $\operatorname{SU}(5)$ with five-fold symmetry of icosian $R_{1}=\frac{1}{2}\left(-\tau+e_{1}+\sigma e_{2}\right)$.

Negatives of these roots constitute the conjugate representation (5*, 10). Similarly the roots belonging to ( $10, \underline{5}$ ) can be written as

$$
\begin{equation*}
R_{1}^{n} B_{m}, \quad-\sigma R_{1}^{n} A_{m} \quad n, m=1,2,3,4,5 \tag{21}
\end{equation*}
$$

the negatives of which represent the roots in ( $1 \underline{0}^{*}, \underline{5}^{*}$ ). It is obvious from (20) and (21) that the roots belonging to the representations $\left(\underline{5}, \underline{0^{*}}\right),(\underline{10}, \underline{5})$ and their conjugates preserve $Z_{5}$ symmetry separately.

It is perhaps more convenient to express the roots of the coset space as products of $R_{1}^{n}$ with the elements of the binary tetrahedral group which can be defined by

$$
\begin{array}{ll}
S_{0}=\frac{1}{2}\left(1+e_{1}+e_{2}+e_{3}\right) & S_{1}=\frac{1}{2}\left(1+e_{1}-e_{2}-e_{3}\right)  \tag{22}\\
S_{2}=\frac{1}{2}\left(1-e_{1}+e_{2}-e_{3}\right) & S_{3}=\frac{1}{2}\left(1-e_{1}-e_{2}+e_{3}\right)
\end{array}
$$

with their conjugates and negatives. Then all the roots of $E_{8}$ can be decomposed as follows:
$(20,1): \pm R_{1}^{n}, \quad \pm \sigma R_{1}^{n} \quad(1,20): \pm R_{1}^{n} e_{3}, \quad \pm \sigma R_{1}^{n} e_{3}$
$\left(\underline{5}, \underline{0^{*}}\right)+\left(\underline{5}^{*}, \underline{10}\right)+(\underline{10}, \underline{5})+\left(\underline{10} \underline{ }^{*}, \underline{5}^{*}\right): \pm R_{1}^{n}\left(S_{\alpha}, \bar{S}_{\beta}, e_{1}, e_{2}\right), \quad \pm \sigma R_{1}^{n}\left(S_{\alpha}, \bar{S}_{\beta}, e_{1}, e_{2}\right)$
$n=1,2,3,4,5 \quad \alpha, \beta=0,1,2,3$.
As we stated at the beginning of this section, $R_{1}$ can be chosen six different ways:

$$
\begin{array}{lll}
R_{1}=\frac{1}{2}\left(-\tau+e_{1}+\sigma e_{2}\right) & R_{2}=\frac{1}{2}\left(-\tau+e_{2}+\sigma e_{3}\right) & R_{3}=\frac{1}{2}\left(-\tau+\sigma e_{1}+e_{3}\right) \\
R_{4}=\frac{1}{2}\left(-\tau+e_{1}-\sigma e_{2}\right) & R_{5}=\frac{1}{2}\left(-\tau+e_{2}-\sigma e_{3}\right) & R_{6}=\frac{1}{2}\left(-\tau-\sigma e_{1}+e_{3}\right) . \tag{24b}
\end{array}
$$

These choices are made so that the triples ( $R_{1}, R_{2}, R_{3}$ ) and ( $R_{4}, R_{5}, R_{6}$ ) have cyclic symmetries in $e_{1}, e_{2}$ and $e_{3}$. It is clear from these discussions that $\mathrm{SU}(5) \times \mathrm{SU}(5)^{\prime}$ can be embedded in $\mathrm{E}_{8}$ in six different ways; in each case a $\mathrm{Z}_{5}$ symmetry is manifest. These six possible decompositions of $\mathrm{E}_{8}$ with respect to $\mathrm{SU}(5) \times \mathrm{SU}(5)^{\prime}$ can be displayed as follows:
$\pm R_{a}^{n}, \pm \sigma R_{a}^{n} \quad \pm R_{a}^{n} e_{3}, \pm \sigma R_{a}^{n} e_{3}$

$$
\begin{align*}
& \left(\underline{5}, \underline{10^{*}}\right)+\left(\underline{5}^{*}, \underline{10}\right)+(\underline{10}, \underline{5})+\left(\underline{10}^{*}, \underline{5}^{*}\right)  \tag{20,1}\\
& \left( \pm R_{a}^{n}, \pm \sigma R_{a}^{n}\right)\left(S_{\alpha}, \bar{S}_{\beta}, e_{1}, e_{2}\right)  \tag{25a}\\
& \left( \pm R_{b}^{n}, \pm \sigma R_{b}^{n}\right)\left(S_{\alpha}, \bar{S}_{\beta}, e_{2}, e_{3}\right) \tag{25b}
\end{align*}
$$

$$
\pm R_{b}^{n}, \pm \sigma R_{b}^{n} \quad \pm R_{b}^{n} e_{1}, \pm \sigma R_{b}^{n} e_{1}
$$

$$
\begin{equation*}
\pm R_{c}^{n}, \pm \sigma R_{c}^{n} \quad \pm R_{c}^{n} e_{2}, \pm \sigma R_{c}^{n} e_{2} \quad\left( \pm R_{c}^{n}, \pm \sigma R_{c}^{n}\right)\left(S_{\alpha}, \bar{S}_{\beta}, e_{3}, e_{1}\right) \tag{25c}
\end{equation*}
$$

where $n=1,2,3,4,5, a=1,4, b=2,5, c=3,6$ and $\alpha, \beta=0,1,3$
So far we have discussed the action of an element of the binary icosahedral group on its elements by left or right multiplication. We can also consider a transformation of an element where a left and right multiplication are combined. Let $P$ and $Q$ be elements of $\langle 5,3,2\rangle$. Then a transformation of the form

$$
\begin{equation*}
Q^{\prime}=( \pm P) Q( \pm \bar{P}) \tag{26}
\end{equation*}
$$

represents the action of an element of the icosahedral group on the elements of $\langle 5,3,2\rangle$. It is also possible to generate a $\mathrm{Z}_{5}$ symmetry in this form. For an illustration of this point let us choose $P=R_{1}=\frac{1}{2}\left(-\tau+e_{1}+\sigma e_{2}\right)$. With this choice of $P$ each $\operatorname{SU}(5)$ root remains invariant without being affected at all by the group operation. However, the


Figure 2. Extended Coxeter-Dynkin diagram of $\mathrm{E}_{8}$ leading to an $\mathrm{SU}(5) \times \operatorname{SU}(5)^{\prime}$ embedding with $Z_{5}$ symmetry.
same action (26) on the roots of $\mathrm{SU}(5)^{\prime}$ manifests itself as a $\mathrm{Z}_{\mathrm{s}}$ symmetry. An extended Coxeter-Dynkin diagram of $E_{8}$, where the $Z_{5}$ symmetries so far discussed are apparent, is given in figure 2 . An $\mathrm{SU}(5)$ orbifold with $\mathrm{Z}_{5}$ symmetry can be used for the construction of six-dimensional string theories [15].

## 4. Embedding of $\operatorname{SU}(3) \times E_{6}$ in $E_{8}$ with $Z_{3}$ symmetry

In this section we decompose $E_{8}$ roots under its maximal subgroup $\mathrm{SU}(3) \times \mathrm{E}_{6}$, which plays an important role in orbifold compactification of the heterotic string. We will choose the simple roots of $E_{8}$ in such a way that the extended Coxeter-Dynkin diagram of $E_{6}$ obtained from that of $E_{8}$ will have a three-fold symmetry of icosians. Such an extended Coxeter-Dynkin diagram of $\mathrm{E}_{6}$ is given in figure 3. Here $Z_{3}$ symmetry is defined in the form of the transformation (26) where $P$ is replaced by $S_{0}=$ $\frac{1}{2}\left(1+e_{1}+e_{2}+e_{3}\right)$, which satisfies $S_{0}^{2}=-\bar{S}_{0}, S_{0}^{3}=-1, S_{0}+\bar{S}_{0}=1$. It is easy to show that an action of $S_{0}$ in the form of (26) rotates $e_{1}, e_{2}, e_{3}$ and $\sigma e_{1}, \sigma e_{2}, \sigma e_{3}$ in cyclic order. We choose the simple roots of $E_{8}$ in such a way that the roots of $E_{6}$ include the elements of the binary tetrahedral group. $\mathrm{SU}(3)$ roots orthogonal to those of $\mathrm{E}_{6}$ in figure 3 are given by $\pm \sigma, \pm \sigma S_{0}, \pm \sigma \bar{S}_{0}$. Each of these roots is invariant under the action of $S_{0}$ in the form of (26). It is clear from figure 3 that $S_{0}$ rotates $B_{1}, B_{2}$ and $B_{3}$ in cyclic order just like it does $e_{1}, e_{2}$ and $e_{3}$. Thus, $\mathrm{Z}_{3}$ symmetry of the extended Coxeter-Dynkin diagram of $E_{6}$ is obtained by repeated application of the element $S_{0}$ of the $E_{8}$ root


Figure 3. Extended Coxeter-Dynkin diagram of $E_{6}$ with icosians $B_{1}=\frac{1}{2}\left(-1-\tau e_{1}-\sigma e_{3}\right)$, $B_{2}=\frac{1}{2}\left(-1-\tau e_{2}-\sigma e_{1}\right), B_{3}=\frac{1}{2}\left(-1-\tau e_{3}-\sigma e_{2}\right), B_{0}=\bar{S}=\left(1-e_{1}-e_{2}-e_{3}\right), e_{1}, e_{2}$ and $e_{3}$.
system. The 72 non-zero roots of $\mathrm{E}_{6}$ can be grouped as 24 triples, each of which is rotated by the action of $S_{0}$. We do not give the details here, but it can be easily shown that while 48 of the roots of $E_{6}$ come from the elements of the binary icosahedral group, the remaining roots are the $\sigma$ multiples of the elements of the icosahedral group lying in the coset space $\mathrm{E}_{8} / \mathrm{SU}(3) \times \mathrm{E}_{6}$.

An interesting algebraic property of the roots $B_{1}, B_{2}$ and $B_{3}$ is the relations [14]:

$$
\begin{align*}
& {\left[B_{i}, B_{j}\right]=-B_{0}-\varepsilon_{i j k} B_{k} \quad B_{i} B_{j} B_{k}=B_{i} S_{0}}  \tag{27}\\
& B_{0}=\bar{S}_{0} \quad i \neq j \neq k=1,2,3 .
\end{align*}
$$

Similar relations can be obtained for the roots cyclically rotated to each other. Let $Q_{1}, Q_{2}$ and $Q_{3}$ be such roots of $\mathrm{E}_{6}$ belonging to $\langle 5,3,2\rangle$. Then we can show that their triple products can be classified as

$$
Q_{1} Q_{2} Q_{3}= \begin{cases}Q_{1} S_{0} & \text { for }\left(Q_{1}, \bar{S}_{0}\right)=0  \tag{28}\\ 1 & \text { for }\left(Q_{1}, \bar{S}_{0}\right)=\frac{1}{2} \\ -1 & \text { for }\left(Q_{1}, \bar{S}_{0}\right)=-\frac{1}{2}\end{cases}
$$

The scalar products must be understood in the 'reduced' form. A cyclic rotation in (28) is also implicit.

Using the notation of (10), we can give an explicit decomposition of the $\dot{E}_{8}$ roots under $\operatorname{SU}(3) \times \mathrm{E}_{6}$ :

$$
\begin{array}{lllll}
(0, & \pm 1 & \pm e_{1} & \pm e_{2} & \left. \pm e_{3}\right) \\
(0, & \pm S_{0} & \pm \bar{S}_{2} & \pm S_{1} & \left. \pm \bar{S}_{1}\right) \\
(0, & \pm \bar{S}_{0} & \pm S_{3} & \pm \bar{S}_{3} & \left. \pm S_{2}\right) \\
\left(A_{1},\right. & \pm \frac{1}{2}\left(e_{3}-e_{1}\right) & \pm \frac{1}{2}\left(1-e_{2}\right) & \pm \frac{1}{2}\left(e_{1}+e_{3}\right) & \left. \pm \frac{1}{2}\left(1+e_{2}\right)\right) \\
\left(A_{2},\right. & \pm \frac{1}{2}\left(e_{1}-e_{2}\right) & \pm \frac{1}{2}\left(1+e_{3}\right) & \pm \frac{1}{2}\left(1-e_{3}\right) & \left. \pm \frac{1}{2}\left(e_{1}+e_{2}\right)\right)  \tag{29b}\\
\left(A_{3},\right. & \pm \frac{1}{2}\left(e_{2}-e_{3}\right) & \pm \frac{1}{2}\left(e_{3}+e_{2}\right) & \pm \frac{1}{2}\left(1+e_{1}\right) & \left. \pm \frac{1}{2}\left(1-e_{1}\right)\right) \\
\left(A_{0},\right. & 0) . & & &
\end{array}
$$

Our notation needs clarification. A bracket ( $a, b c d e$ ) represents four brackets of the form $(a, b),(a, c),(a, d),(a, e)$. In this unusual notation $(0, \pm 1)= \pm \sigma,\left(0, \pm S_{0}\right)= \pm \sigma S_{0}$, $\left(0, \pm \bar{S}_{0}\right)= \pm \sigma \bar{S}_{0}$ are the roots of $\mathrm{SU}(3) . \mathrm{E}_{6}$ roots are given in the same column as $\mathrm{SU}(3)$ roots in (29) which can be also written as
$A_{0} \quad A_{1} \pm \frac{1}{2} \sigma\left(e_{3}-e_{1}\right) \quad A_{2} \pm \frac{1}{2} \sigma\left(e_{1}-e_{2}\right) \quad A_{3} \pm \frac{1}{2} \sigma\left(e_{2}-e_{3}\right)$.
Since $S_{0}$ is an element of $A_{0}$ it is left unchanged by the action of $S_{0}$. The other sets of elements are rotated into each other in cyclic order since $A_{1}, A_{2}$ and $A_{3}$ are rotated into each other by $S_{0}$ in the form of (26). It is quite obvious that the remaining elements belong to the coset space $\mathrm{E}_{8} / \mathrm{SU}(3) \times \mathrm{E}_{6}$ and are separately rotated into each other by $S_{0}$.

The roots of $E_{8}$ in (29) are organised such that each column is invariant under right multiplication by $S_{0}$, in contrast to the transformation defined by (26). While each column in (29a) represents the roots of $\mathrm{SU}(3)$, the column below displays the roots of $\mathrm{E}_{6}$ provided ( $A_{0}, 0$ ) is also added. Therefore, the root system given in (29) shows that $\mathrm{SU}(3) \times \mathrm{E}_{6}$ can be embedded in $\mathrm{E}_{8}$ in four different ways with an obvious $S_{0}$ invariance by right multiplication. However, an action of $S_{0}$ in the sense of (26)
will allow only one decomposition. Similar decompositions of $\mathrm{E}_{8}$ roots under $\mathrm{SU}(3) \times$ $\mathrm{E}_{6}$ can be made, replacing $S_{0}$ by any icosian which satisfies $q^{3}= \pm 1$.

String phenomenologists sometimes prefer an $\mathrm{SU}(3)^{3}$ orbifold with $\mathrm{Z}_{3}$ symmetry [3]. The simple roots of $E_{8}$ can be arranged so that a $Z_{3}$ symmetry for each $\operatorname{SU}(3)^{4} \subset$ $\mathrm{SU}(3) \times \mathrm{E}_{6} \subset \mathrm{E}_{8}$ can be realised. For this purpose let us choose a root $R=$ $\frac{1}{2}\left(-1+\tau e_{1}+\sigma e_{3}\right)$. Then it can be shown that each of the following sets of roots:

$$
\begin{array}{lr}
\left( \pm R, \pm R^{2}, \pm 1\right) & \left( \pm R e_{2}, \pm R^{2} e_{2}, \pm e_{2}\right) \\
\left( \pm \sigma R, \pm \sigma R^{2}, \pm \sigma\right) & \left( \pm \sigma R e_{2}, \pm \sigma R^{2} e_{2}, \pm \sigma e_{2}\right) \tag{31}
\end{array}
$$

represents one $\mathrm{SU}(3)$. Since $R^{3}=1$, each $\mathrm{SU}(3)$ has a $Z_{3}$ symmetry. The $Z_{3}$ symmetry discussed in this section can also be represented by the octonionic roots of $E_{8}$, which will be discussed in a separate publication [16].

## 5. Discussion and conclusion

Several aspects of the representation of the $\mathrm{E}_{8}$ lattice with icosians differ from the octonionic representation. The two descriptions of the $\mathrm{E}_{8}$ lattice can be contrasted as follows.
(i) The octonionic root system obeys the usual scalar product defined by (7) and forms a closed non-associative discrete algebra of order 240 , only 24 of which satisfy the group property of the binary tetrahedral group.
(ii) In the case of icosians, the order of the group structure is extended to the binary icosahdral group of order 120 but the whole set of 240 roots do not close under multiplication since a multiplication of the form $(\sigma q)(\sigma q)=\sigma^{2} q^{2}=q+\sigma q$ produces lattice vectors of higher norms. Icosians represent the $\mathrm{E}_{8}$ lattice only with the 'reduced' scalar product.
(iii) Octonionic roots yield natural Abelian symmetries $Z_{6}, Z_{4}, Z_{3}$ and $Z_{2}$ of the $\mathrm{E}_{8}$ lattice with an interesting manifestation of the triality of the extended CoxeterDynkin diagram of $\mathrm{E}_{6}$ [14].
(iv) With icosians, while preserving the three-fold symmetry of the extended Coxeter-Dynkin diagram of $\mathrm{E}_{6}$, one can naturally extend the Abelian symmetries to $Z_{10}, Z_{6}, Z_{5}, Z_{4}, Z_{3}$ and $Z_{2}$ of the root system of $E_{8}$. To be more specific, the maximal subgroup $\mathrm{SU}(5) \times \operatorname{SU}(5)^{\prime}$ can be embedded in $\mathrm{E}_{8}$ with a $\mathrm{Z}_{5}$ symmetry invariance, which is not possible in the octonionic representation of the $\mathrm{E}_{8}$ lattice.

Another amusing observation is the possibility of describing the $\mathrm{E}_{8} \times \mathrm{E}_{8}^{\prime}$ root system by a simple extension of the root system given by (10). If we multiply the icosianic roots in (10) by the octonionic imaginary unit $e_{7}$ we obtain an independent root system of 240 elements described by the octonionic units $e_{4}, e_{5}, e_{6}, e_{7}$ and their $\sigma$ multiples. Then one can show that the 120 elements of the binary icosahedral group in $\mathrm{E}_{8}$ and the corresponding 120 octonionic elements in $\mathrm{E}_{8}^{\prime}$ form a closed algebra, presumably a subset of a larger algebra. The products of any two elements of the $E_{8}^{\prime}$ lattice will yield a lattice element of $E_{8}$. Another possible description of the $\mathrm{E}_{8} \times \mathrm{E}_{8}^{\prime}$ lattice can be made by multiplying the octonionic roots of $\mathrm{E}_{8}$ by $\sigma$. This second construction is totally different from the previous one in most respects.

The first construction of $\mathrm{E}_{8} \times \mathrm{E}_{8}^{\prime}$ could be attractive in view of the fact that an unbroken $\mathrm{E}_{8}^{\prime}$ in the heterotic string could be attributed to its pure octonionic structure.

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